# Dynamical properties of neuron models-nodal and collective behaviours.

Indranil Ghosh

School of Mathematical and Computational Sciences Massey University, Palmerston North, New Zealand

August 14, 2024



▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - 釣��

#### Neurons as Dynamical units

- ▶ Neurons represent the fundamental dynamical units of the nervous system
- The dynamics of neurons, like firing of action potentials, can be modeled as simple dynamical systems like ODEs or maps

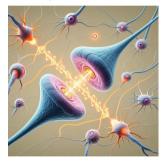


Figure: Two neurons connected by a synapse. (Powered by DALL-E 3)

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ の 0 0

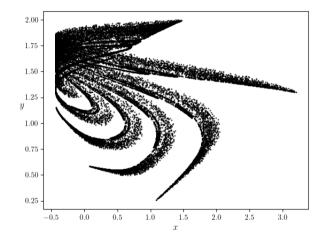
## Chialvo Map (Chialvo, 1995)

The two-dimensional neuron map is given by

$$x_{n+1} = x_n^2 e^{(y_n - x_n)} + k_0,$$
  
 $y_{n+1} = ay_n - bx_n + c.$ 

- The state variables x and y represent the activation variable and recovery-like variable,
- $\blacktriangleright$  a, b, c and  $k_0$  are the system parameters,
- ▶ *a* < 1 is the time constant of recovery,
- $\blacktriangleright$  *b* < 1 represents the activation dependence of the recovery process,
- c denotes the offset, and
- $\triangleright$   $k_0$  is the time-independent additive perturbation.

#### A Typical Phase Portrait



◆□▶ ◆圖▶ ◆臣▶ ◆臣▶ 臣 のへで

#### Electromagnetic flux

We describe the effects of electromagnetic flux on the system of neurons with **memristors**. The induction current due to electromagnetic flux is given by

$$rac{dq(\phi)}{dt} = rac{dq(\phi)}{d\phi} rac{d\phi}{dt} = M(\phi) rac{d\phi}{dt} = kM(\phi)x.$$

- $\blacktriangleright$   $\phi$ : electromagnetic flux across the neuron membranes,
- k: electromagnetic flux coupling strength, &
- $M(\phi)$ : memconductance of electromagnetic flux controlled memristor. We consider the following memconductance function:

$$M(\phi) = \alpha + 3\beta\phi^2.$$

# Improved Chialvo map under electromagnetic flux (Muni, Fatoyinbo, & Ghosh, 2022)

International Journal of Bifurcation and Chaos, Vol. 32, No. 9 (2022) 2230020 (26 pages) © World Scientific Publishing Company DOI: 10.1142/S0218127422300208

#### Dynamical Effects of Electromagnetic Flux on Chialvo Neuron Map: Nodal and Network Behaviors

Sishu Shankar Muni<sup>\*,†,†</sup>, Hammed Olawale Fatoyinbo<sup>†,§</sup> and Indranil Ghosh<sup>†,¶</sup> "Department of Physical Sciences, Indian Institute of Science and Educational Research Kolkata, Campus Road, Mohanpur, West Bengal 741246, India <sup>†</sup>School of Mathematical and Computational Sciences, Massey University, Colombo Road, Palmerston North, 4410, New Zealand <sup>†</sup>sishul 1729@iierschol.a.c.in <sup>†</sup>s.muni@massey.ac.nz <sup>§</sup>h.fatojinbo@massey.ac.nz <sup>§</sup>h.fatojinbo@massey.ac.nz

Received January 11, 2022; Revised May 4, 2022

Improved Chialvo map under electromagnetic flux (Muni, Fatoyinbo, & Ghosh, 2022)

Under the action of electromagnetic flux, the system of Chialvo map is improved to the following map:

$$x_{n+1} = x_n^2 e^{(y_n - x_n)} + k_0 + k x_n M(\phi_n),$$
  

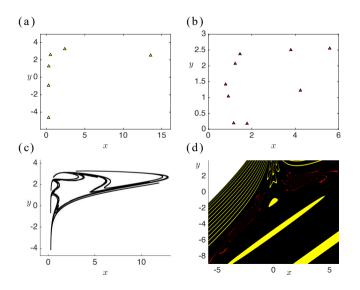
$$y_{n+1} = a y_n - b x_n + c,$$
  

$$\phi_{n+1} = k_1 x_n - k_2 \phi_n,$$

making the system a three-dimensional smooth map. The new variables  $\alpha$ ,  $\beta$ ,  $k_1$ ,  $k_2$  represent the electromagnetic flux parameters.

◆□▶ ◆□▶ ◆目▶ ◆目▶ ●目 ● のへの

## Multistability



◆□ ▶ ◆□ ▶ ◆ □ ▶ ◆ □ ▶ ● □ ● ● ● ●

#### Bifurcation structures and antimonotonicity

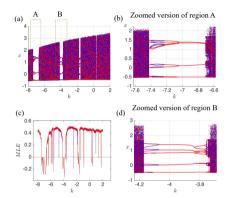


Figure: Bifurcation diagram of x with respect to k in panel (a). A maximal Lyapunov exponent diagram is shown in panel (b).

#### Bifurcation structures and antimonotonicity

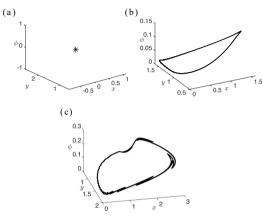


Figure: In (a) a stable fixed point is shown in the  $x - y - \phi$  phase space for a = 0.838. After a supercritical Neimark-Sacker bifurcation, an attracting closed invariant curve is born as shown in (b) at a = 0.841. A chaotic attractor is then formed when a is increased to 0.88.

#### Numerical bifurcation analysis

Codimension-1			
Saddle-node (fold) bifur-	LP	Neimerk-Sacker bifurca-	NS
cation		tion	
Period-doubling (flip) bi-	PD		
furcation			
Codimension-2			
Cusp	CP	Chenciner	СН
Generalized flip	GPD	Fold-Flip	LPPD
Flip-Neimark-Sacker	PDNS	Fold-Neimark-Sacker	LPNS
1:1 resonance	R1	1:2 resonance	R2
1:3 resonance	R3	1:4 resonance	R4

#### Table: Abbreviations of codimension-1 and codimension-2 bifurcations

#### Numerical bifurcation analysis

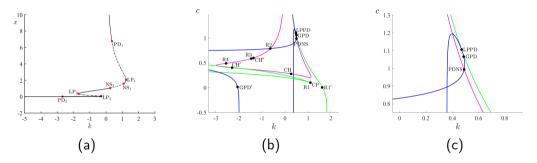


Figure: (a) Codimension-1 bifurcation diagram with k as bifurcation parameter. (b) Codimesion-2 bifurcation diagram in (k, c)-parameter plane. (c) Zoomed version of (b)

#### Bursting and spiking features Tonic spiking Regular spiking (a) k = 0.6, b = 0.8(b) k = 0.6, b = 0.95 xx5 0 0 100 200 300 0 100 200 300 0 Periodic bursting (c) Phasic bursting k = 0.6, b = 0.5 (d) k = 0.4, b = 0.110 $_x$ 5 . x5 0 0 100 200 0 0 300 50 100 t t (e) Chaotic bursting k = -0.5, b = 0.1 $x^{10}$ ' 5 0 0 100 200 300 ж

## Ring-star network for multiple neurons

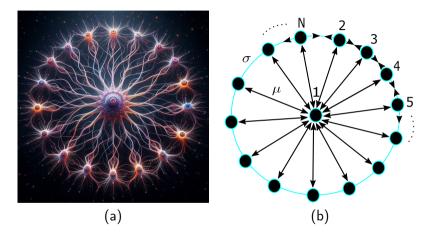


Figure: (a) Ensemble of connected neurons. (Powered by DALL-E 3). (b) Ring-star network of Chialvo neuron system.

#### Ring-star network for multiple neurons

The mathematical model for the ring-star connected Chialvo neuron map under electromagnetic flux is defined as:

$$\begin{aligned} x_m(n+1) &= x_m(n)^2 e^{y_m(n) - x_m(n)} + k_0 + k x_m(n) \mathcal{M}(\phi_m(n)) \\ &+ \mu(x_m(n) - x_1(n)) + \frac{\sigma}{2R} \sum_{i=m-R}^{m+R} (x_i(n) - x_m(n)), \\ y_m(n+1) &= a y_m(n) - b x_m(n) + c, \\ \phi_m(n+1) &= k_1 x_m(n) - k_2 \phi_m(n), \end{aligned}$$

#### Ring-star network for multiple neurons

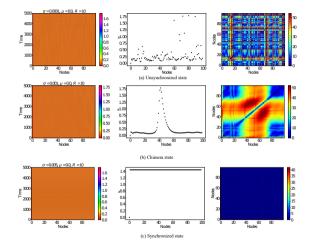
▶ The central node is further defined as

$$\begin{aligned} x_1(n+1) &= x_1(n)^2 e^{(y_1(n)-x_1(n))} + k_0 + k x_1(n) M(\phi_1(n)) \\ &+ \mu \sum_{i=1}^N (x_i(n)-x_1(n)), \\ y_1(n+1) &= a y_1(n) - b x_1(n) + c, \\ \phi_1(n+1) &= k_1 x_1(n) - k_2 \phi_1(n), \end{aligned}$$

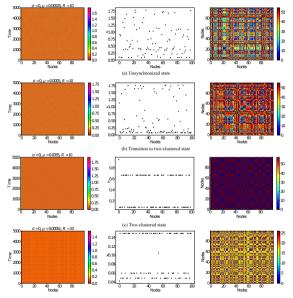
having the following boundary conditions:

$$\begin{aligned} x_{m+N}(n) &= x_m(n), \qquad (1) \\ y_{m+N}(n) &= y_m(n), \\ \phi_{m+N}(n) &= \phi_m(n). \end{aligned}$$

#### Simulations



#### Simulations



(d) Three-clustered state

Nonlinear Dyn (2023) 111:17499–17518 https://doi.org/10.1007/s11071-023-08717-y



▲□▶ ▲□▶ ▲□▶ ▲□▶ □ の 0 0

ORIGINAL PAPER

## On the analysis of a heterogeneous coupled network of memristive Chialvo neurons

Indranil Ghosh · Sishu Shankar Muni · Hammed Olawale Fatoyinbo

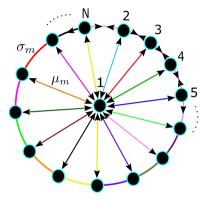


Figure: The star and ring coupling strengths are denoted by  $\mu_m$  and  $\sigma_m$  for each node m = 1, ..., N respectively. Different colors in the ring-star topology signify a range of heterogeneous values of  $\mu_m$  and  $\sigma_m$ .

- We introduce heterogeneities to the coupling strengths  $\sigma_m(n)$  and  $\mu_m(n)$  both in space and time.
- In space, the heterogeneities are realized following the application of a noise source with a uniform distribution given by

$$\sigma_m(n) = \sigma_0 + D_\sigma \xi_\sigma^{m,n},\tag{2}$$

$$\mu_m(n) = \mu_0 + D_\mu \xi_\mu^{m,n}, \tag{3}$$

- ▶ Here  $\sigma_0$  and  $\mu_0$  are the mean values of the coupling strengths  $\mu_m$  and  $\sigma_m$  respectively.
- We keep  $\sigma_0 \in [-0.01, 0.01]$  and  $\mu_0 \in [-0.001, 0.001]$ .
- The noise sources  $\xi_{\sigma}$  and  $\xi_{\mu}$  for the corresponding coupling strengths are real numbers randomly sampled from the uniform distribution [-0.001, 0.001].
- Finally, the D's refer to the "noise intensity" which we restrict in the range [0, 0.1].

- Heterogeneity in time is introduced by considering the network having time-varying links depending on the two coupling probabilities  $P_{\mu}$  and  $P_{\sigma}$ , which govern the update of the coupling topology with each iteration n.
- The probability with which the central node is connected to all the peripheral nodes at a particular n is denoted by P<sub>µ</sub>.
- Likewise, the probability with which the peripheral nodes are connected to their R neighboring nodes is given by P<sub>σ</sub>.
- We employ three metrics to analyse our model: (1) cross-correlation coefficient,
   (2) synchronization error, and (3) Sample entropy

#### Quantitative metrics

 $\blacktriangleright$  The general definition of the cross-coefficient denoted by  $\Gamma_{i,m}$  is given by

$$\Gamma_{i,m} = \frac{\langle \tilde{x}_i(n)\tilde{x}_m(n)\rangle}{\sqrt{\langle (\tilde{x}_i(n))^2 \rangle \langle (\tilde{x}_m(n))^2 \rangle}}.$$
(4)

The averaged cross-correlation coefficient over all the units of the network is given by,

$$\Gamma = \frac{1}{N-1} \sum_{m=1, m \neq i}^{N} \Gamma_{i,m}.$$
(5)

- We use Γ<sub>2,m</sub>, denoting the degree of correlation between the first peripheral node of the ring-star network and all the other nodes, including the central node.
- The average is calculated over time with transient dynamics removed and  $\tilde{x}(n) = x(n) \langle x(n) \rangle$ .

> The averaged synchronization-error for the nodes in a system is given by

$$E = \frac{1}{N-1} \sum_{m=1, m \neq 2}^{N} \langle |x_2(n) - x_m(n)| \rangle,$$
 (6)

• We again consider node number N = 2 as the baseline.

#### Simulations

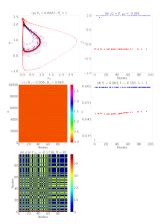


Figure: Coherent and solitary nodes giving rise to a two-clustered state.

#### simulations

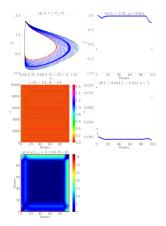


Figure: Mostly synced in the coherent domain with two solitary nodes.

#### Simulations

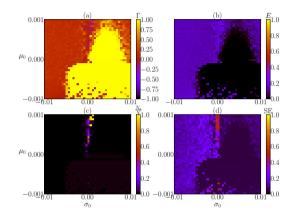


Figure: An almost definitive bifurcation boundary is observed. Solitary nodes appear around  $\sigma_0 \sim 0$  and  $\mu_0 > 0$ .

#### Simulations

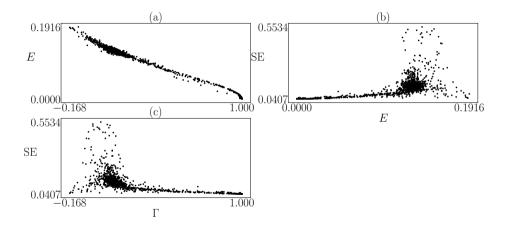


Figure: Comparison plots for the various measures. Figures (a) and (c) show an inverse trend whereas figure (b) shows a proportional trend.

## A bit of a digression! (Ghosh, Nair, Fatoyinbo, & Muni, 2024)

Eur. Phys. J. Plus (2024) 139:545 https://doi.org/10.1140/epjp/s13360-024-05363-0

Regular Article

#### THE EUROPEAN PHYSICAL JOURNAL PLUS



## Dynamical properties of a small heterogeneous chain network of neurons in discrete time

#### Indranil Ghosh<sup>1,a</sup>, Anjana S. Nair<sup>2,b</sup>, Hammed Olawale Fatoyinbo<sup>3,4,c</sup>, Sishu Shankar Muni<sup>2,d</sup>

<sup>1</sup> School of Mathematical and Computational Sciences, Massey University, Colombo Road, Palmerston North 4410, New Zealand

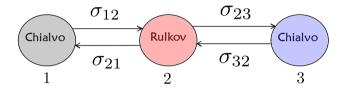
<sup>2</sup> School of Digital Sciences, Digital University Kerala, Technopark Phase-IV Campus, Mangalapuram, Kerala 695317, India

<sup>3</sup> EpiCentre, School of Veterinary Science, Massey University, Palmerston North 4410, New Zealand

<sup>4</sup> Department of Mathematical Sciences, Auckland University of Technology, Auckland 1010, New Zealand

Received: 25 March 2024 / Accepted: 13 June 2024 © The Author(s) 2024

A bit of a digression! (Ghosh, Nair, Fatoyinbo, & Muni, 2024)



▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQ@

Figure: A heterogeneous network of a tri-oscillator chain composed of end nodes (Chialvo neuron map) and central node (Rulkov neuron map).

Higher-order smallest ring-star network (Nair, Ghosh, Fatoyinbo, & Muni, 2024)

# On the higher-order smallest ring-star network of Chialvo neurons under diffusive couplings



#### AFFILIATIONS

 School of Digital Sciences, Digital University Kerala, Technopark Phase-IV campus, Mangalapuram 695317, Kerala, India
 School of Mathematical and Computational Sciences, Massey University, Colombo Road, Palmerston North 4410, New Zealand
 Department of Mathematical Sciences, School of Engineering, Computer and Mathematical Sciences, Auckland University of Technology, Auckland 1142, New Zealand

a)Author to whom correspondence should be addressed: i.ghosh@massey.ac.nz

# Higher-order smallest ring-star network (Nair, Ghosh, Fatoyinbo, & Muni, 2024)

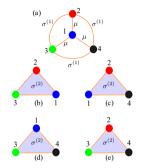


Figure: In panel (a) the coupling strength within the star configuration is denoted by  $\mu$ , while the coupling strength within the ring is denoted by  $\sigma^{(1)}$ . Moreover, the coupling strength originated from the higher-order interactions is represented by  $\sigma^{(2)}$ , as indicated by the triplets in panels (b) $\rightarrow$ (e).

#### Model

This system in compact form is written as

$$\begin{aligned} x_{p}(n+1) &= x_{p}(n)^{2} e^{(y_{p}(n) - x_{p}(n))} + k_{0} + \mu(x_{1}(n) - x_{p}(n)) \\ &+ \sigma^{(1)} \sum_{i=2}^{4} (x_{i}(n) - x_{p}(n)) \\ &+ \sigma^{(2)} \sum_{i=1}^{4} \sum_{\substack{j=i+1\\ i \neq p\\ j \neq p}}^{4} (x_{i}(n) + x_{j}(n) - 2x_{p}(n)), \end{aligned}$$
(7)  
$$y_{p}(n+1) &= ay_{p}(n) - bx_{p}(n) + c, \end{aligned}$$
(8)

<日 > < 団 > < 団 > < 団 > < 団 > < 団 > < 団 > < 団 > < 団 < 〇 < 〇</p>

Model

 $\mathsf{and}$ 

$$x_{1}(n+1) = x_{1}(n)^{2} e^{(y_{1}(n)-x_{1}(n))} + k_{0} + \mu \sum_{i=2}^{4} (x_{i}(n) - x_{1}(n)) + \sigma^{(2)} \sum_{i=2}^{4} \sum_{j=i+1}^{4} (x_{i}(n) + x_{j}(n) - 2x_{1}(n)), \qquad (9) y_{1}(n+1) = ay_{1}(n) - bx_{1}(n) + c. \qquad (10)$$

#### Simulations

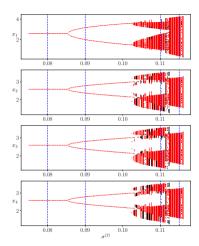


Figure: Bifurcation plot of each node against the coupling strength  $\sigma^{(2)}$  once simulated forward (points colored black) and once backward (points colored red).

#### Simulations

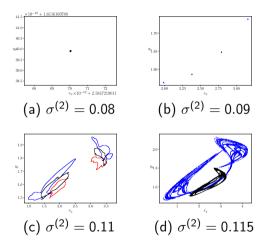


Figure: Typical phase portraits. (a) fixed point, (b) period-doubling, (c) a disjoint cyclic quasiperiodic closed invariant curve, (d) chaos.

0-1 test

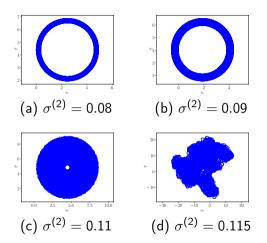


Figure: Signal plots. (a) highly bounded trajectory, (b) slightly less bounded trajectory, (c) between bounded and diffusive, and (d) diffusive random walk corresponding to a Brownian motion with zero drift.

#### Metrics

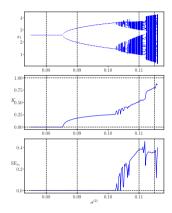


Figure: A bifurcation plot of the first node with  $\sigma^{(2)}$  as the main bifurcation parameter. The corresponding value of K from the chaos test and SE for complexity are shown.

#### Denatured Morris-Lecar model (Fatoyinbo et al., 2022)

2022 International Conference on Decision Aid Sciences and Applications (DASA)

#### Numerical Bifurcation Analysis of Improved Denatured Morris-Lecar Neuron Model

Hammed Olawale Fatoyinbo School of Mathematical and Computational Sciences Massey University Palmerston North, New Zealand h.fatoyinbo@massey.ac.nz ORCID: 0000-0002-6036-2957

Indranil Ghosh School of Mathematical and Computational Sciences Massey University Palmerston North, New Zealand i.ghosh@massey.ac.nz Sishu Shankar Muni School of Mathematical and Computational Sciences Massey University Palmerston North, New Zealand s.muni@massey.ac.nz ORCID: 0000-001-9545-8345

Ibrahim Olatunji Sarumi Department of Mathematics and Statistics King Fahd University of Petroleum and Minerals Dhahran, Saudi Arabia isarumi@kfupm.edu.sa

Afeez Abidemi Department of Mathematical Sciences Federal University of Technology Akure, Nigeria aabidemi@futa.edu.ng ORCID: 0000-0003-1960-0658

◆□ > ◆□ > ◆豆 > ◆豆 > ̄豆 = ∽ 9 < @ ->

#### Denatured Morris-Lecar model (Fatoyinbo et al., 2022)

The dentaured Morris-Lecar model proposed in the book<sup>1</sup> consists of two nonlinearly coupled ODEs

$$\dot{x} = x^2(1-x) - y + I,$$
 (11)

$$\dot{y} = Ae^{\alpha x} - \gamma y. \tag{12}$$

Here x is the action potential, y is again the recovery variable and I is the external current. The other parameters are all positive constants.

<sup>&</sup>lt;sup>1</sup>D. Schaeffer and J. Cain, Ordinary differential equations: Basics and beyond (Springer, 2018) 🚊 🔗 < 📀

#### New Model

$$\dot{x} = x^2(1-x) - y + I,$$
 (13)

$$\dot{y} = Ae^{\alpha x} - \gamma y, \tag{14}$$

$$\dot{I} = \varepsilon (I'(x) - I), \tag{15}$$

(17)

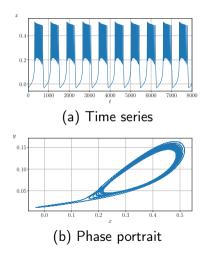
where

$$I'(x) = \frac{1}{60} \left[ 1 + \tanh\left(\frac{0.05 - x}{0.001}\right) \right]$$
(16)

is the smoothed-out version of the step function given by

$$H(x) = egin{cases} rac{1}{30}, & x < 0.05, \ 0, & x > 0.05. \end{cases}$$

#### New Model



◆□ → ◆□ → ◆三 → ◆三 → ○ ● ● ● ●

#### The End

Thank you! Questions?

