BIFURCATION STRUCTURE WITHIN ROBUST CHAOS FOR PIECEWISE-LINEAR MAPS. I. Ghosh, R. McLachlan, D.J.W. Simpson School of Mathematical and Computational Sciences, Massey University, Palmerston North, New Zealand

Introduction

Here I use geometry to explain robust chaotic dynamics in piecewise-linear (PWL) maps. PWL maps are used for modeling systems with switches, thresholds, and other abrupt events. We study the two-dimensional bordercollision normal form:

with variables $x, y \in \mathbb{R}$, and parameter vector $\xi = (\tau_L, \delta_L, \tau_R, \delta_R) \in \mathbb{R}^4$.

A renormalisation operator

Although the second iterate f_{ξ}^2 has four pieces, for many values of ξ only two of these are relevant:

$$f_{\xi}^{2}(x,y) = \begin{cases} \begin{bmatrix} \tau_{L}\tau_{R} - \delta_{L} & \tau_{R} \\ -\delta_{R}\tau_{L} & -\delta_{R} \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} + \begin{bmatrix} \tau_{R} + 1 \\ -\delta_{R} \end{bmatrix} \\ \begin{bmatrix} \tau_{R}^{2} - \delta_{R} & \tau_{R} \\ -\delta_{R}\tau_{R} & -\delta_{R} \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} + \begin{bmatrix} \tau_{R} + 1 \\ -\delta_{R} \end{bmatrix} ,$$

Then f_{ξ}^2 is equivalent to $f_{q(\xi)}$ under a change of coordinates, where g is the renormalisation operator defined by

$$\begin{aligned} \tilde{\tau}_L &= \tau_R^2 - 2\delta_R, \\ \tilde{\delta}_L &= \delta_R^2, \\ \tilde{\tau}_R &= \tau_L \tau_R - \delta_L - \delta_R, \\ \tilde{\delta}_R &= \delta_L \delta_R. \end{aligned}$$

The classical robust chaos parameter region

Robust chaos refers to the absence of periodic windows. We consider the parameter region

 $\Phi = \{ \xi \in \mathbb{R}^4 \mid \tau_L > |\delta_L + 1|, \ \tau_R < |\delta_R + 1| \},\$ where $f_{\mathcal{E}}$ has two saddle fixed points X and Y, see Fig. 1. Within Φ $\Phi_{\rm BYG} = \{\xi \in \Phi \,|\, \delta_L > 0, \delta_R > 0, \phi_4(\xi) > 0\},\$ where $\phi_4(\xi) = \delta_R - (\tau_R + \delta_L + \delta_R - (1 + \tau_R)\lambda_L^u)\lambda_L^u$, is the classical robust chaos parameter region of [1]. For all $n \ge 0$, define $\zeta_n(\xi) = \phi_4(g^n(\xi)),$

and

$$\mathcal{R}_n = \{\xi \in \Phi \mid \zeta_n(\xi) > 0, \zeta_{n+1}(\xi) \le 0\}$$

see Fig. 2. Our main result is that for all $n \ge 0$ and ξ
has a chaotic attractor with 2^n connected components [2]
was obtained by using the renormalisation operator and
unknown bifurcation structure within the robust chaos particular

 $\leq 0,$

 $\geq 0,$

 $, x \leq 0,$

 $x \ge 0.$

 $\xi \in \mathcal{R}_n$ the map f_{ξ} [2], e.g. Fig. 3. This nd reveals previously arameter region.



Fig. 1: A phase portrait with $\xi \in \mathcal{R}_0 \subset \Phi_{BYG}$, showing the initial line segments of the stable and unstable manifolds of the fixed points X and Y, as well as an additional segment of the stable manifold of Y. The attractor (computed numerically) is shown in black.





Invariant expanding cones

Chaos in Φ_{BYG} can be proved by constructing an invariant expanding cone, Fig. 4, in tangent space [3]. We have extended this to Φ ; Figs. 5 and 6 show parameter values for which we have been able to explicitly construct a trapping region and a cone. For any $\xi \in \Phi_{trap} \cap \Phi_{cone}$, f_{ξ} has a chaotic attractor [4].



Fig. 4: A sketch of an invariant expanding cone C and its image $AC = \{Av | v \in C\}, \text{ given } A \in \mathbb{R}^{2 \times 2}.$

Fig. 2: Two-dimensional cross-sections of the parameter regions \mathcal{R}_n , where \mathcal{R}_n is visible for all $n = 0, 1, \ldots, 4$. The region $\Phi_{\rm BYG}$ is bounded by $\tau_L = \delta_L + 1$, $\tau_R = -(\delta_R + 1), \, \zeta_0 = 0, \, \delta_L > 0,$ and $\delta_R > 0$.



 Φ_{trap} , showing curves $\phi_i(\xi) = 0$, for $\tau_L = 1.6$ and $\tau_R = -1.5$.

Extension to higher dimensions

In the N-dimensional setting, suppose the fixed point Y has exactly one unstable eigenvalue $\lambda_L^1 > 1$ and the fixed point X has exactly one unstable eigenvalue $\lambda_R^1 < -1$. We have been able to construct an N-dimensional trapping region in an open region of parameter space, see Fig. 7.



Fig. 7: Our trapping region construction is valid when the absolute values of the stable eigenvalues are all less than the indicated value of r.

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[1] S. Banerjee, J.A. Yorke, 80(14):3049–3052, 1998. [2] I. Ghosh, and D.J.W. Sir 32(12):2250181, 2022. [3] P.A. Glendinning, and D.J.W. 41(7):3367–3387, 2021. [4] I. Ghosh, R. McLachlan, and tion for robust chaos in two-dimensional piecewise-linear maps. In preparation.

for $\tau_L = 1.6$ and $\tau_R = -1.5$.

ferences
and C. Grebogi. Phys. Rev. Lett.,
mpson. Int. J. Bifurcation Chaos,
Simpson. Discrete Contin. Dyn. Syst.,
D.J.W. Simpson. A generalised construc-